Analysis of Hyperelastic Materials with MECHANICA – Theory and Application Examples –

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Part 1

Theoretic Background Information
Review of Hooke’s law for linear elastic materials (1)

- Fundamental equation, well known to all engineers, is:
  \[ \sigma = E \cdot \varepsilon \]
  In this equation, the proportionality constant \( E \) between strain and stress is the “Modulus of Elasticity” of the material.

- Hooke’s law is not as simple as it looks like above: This equation is just valid for the special case of uniaxial tension and in the direction of this tension!

- In order to cover three-dimensional stress and strain states, in a first step we solve this equation for \( \varepsilon \) and just look for the first principal strain:
  \[ \varepsilon_1 = \frac{1}{E} \cdot \sigma_1 \]

- Now, we add on the right side of this equation the missing terms from the two lateral principal stresses \( \sigma_2 \) and \( \sigma_3 \). Compared to \( \sigma_1 \), these lateral stresses influence the first principal strain \( \varepsilon_1 \) much less: So, they are multiplied with a “proportionality constant” \( \leq 0.5 \), known as the Poisson’s ratio \( \nu \):
  \[ \varepsilon_1 = \frac{1}{E} \cdot \{\sigma_1 - \nu(\sigma_2 + \sigma_3)\} \]
Review of Hooke’s law for linear elastic materials (2)

- If we do the same in the other two orthogonal principal directions, we obtain the general formulation of Hooke’s law:
  \[
  \varepsilon_1 = \frac{1}{E} \cdot \{\sigma_1 - \nu(\sigma_2 + \sigma_3)\}
  \]
  \[
  \varepsilon_2 = \frac{1}{E} \cdot \{\sigma_2 - \nu(\sigma_1 + \sigma_3)\}
  \]
  \[
  \varepsilon_3 = \frac{1}{E} \cdot \{\sigma_3 - \nu(\sigma_1 + \sigma_2)\}
  \]

- The limits of the Poisson ratio \(\nu\) are:
  - \(\nu=0\): no influence of lateral stresses to the strain (no lateral contraction)
  - \(\nu=0.5\): incompressible material, means there is no volume change under loads (of course there is usually a big change in shape under loads!)
  - \(\nu=0.2…0.3\): typical values for linear elastic material like ceramic & metal
The strain energy density of linear elastic materials (1)

- When loading and unloading a linear elastic material, we “drive” along the same straight line in the stress-strain characteristic curve:

\[ \sigma \]

\[ E \]

\[ \varepsilon \]

- The strain energy density \( W \) of such a material is expressed as the half value of the double dot product of stress tensor \( S \) and strain tensor \( E \):

\[ W = \frac{1}{2} S \cdot \cdot E \]

- To explain it more simply for listeners who are not familiar with tensor operations, let’s have a look at a simple spring: Every engineer knows its spring energy is

\[ E_{spring} = \frac{1}{2} K\Delta l^2 \]

with \( K \) = spring stiffness and \( \Delta l \) = spring elongation.
The strain energy density of linear elastic materials (2)

- If our spring is a simple tension rod, its spring constant becomes $K = \frac{EA}{l}$ (A=cross section, l=rod length), so we obtain for the spring energy with $\varepsilon = \frac{\Delta l}{l}$

$$E_{spring} = \frac{1}{2} \frac{EA}{l} (l\varepsilon)^2 = \frac{1}{2} EA l \varepsilon^2$$

- The strain energy density $W$ now is the spring energy within each unit volume of the spring. Since for the simple tension rod we have $V = Al$, we obtain:

$$W = \frac{1}{2} E \varepsilon^2$$

- With $\sigma = E \varepsilon$ we can conclude for the strain energy density of uniaxially loaded linear elastic material:

$$W = \frac{1}{2} \sigma \varepsilon$$

- This is exactly the area below the stress-strain curve:

![Diagram](image-url)
Hyperelastic material (1)

Hyperelastic and linear elastic material:

- A hyperelastic material is still an elastic material, that means it returns to it's original shape after the forces have been removed.
- Hyperelastic material also is Cauchy-elastic, which means that the stress is determined by the current state of deformation, and not the path or history of deformation.
- The difference to linear elastic Material is, that in hyperelastic material the stress-strain relationship derives from a strain energy density function, and not a constant factor.
- This definition says nothing about the Poisson's ratio or the amount of deformation that a material will undergo under loading.
- However, often elastomers are modeled as hyperelastic. Hyperelasticity may also be used to describe biological materials, like tissue.

Hyperelastic material behavior

\[ \sigma \quad \varepsilon \]

\[ E_0 \]
Hyperelastic material (2)

Elastomers are often modeled as hyperelastic.

- Elastomers (like rubber) typically have large strains (often some 100 %) at small loads (means a very low modulus of elasticity, for example just 10 MPa). The material is nearly incompressible, so the Poisson’s ratio is very close to 0.5.

- Their loading and unloading stress-strain curve is not the same, depending on different influence factors (time, static or dynamic loading, frequency…). This viscous behavior is ignored if the hyperelastic material model is used for description.

Elastomer material behavior

Typical measurements are often performed for example until 300 % strain.

Linear elastic material (like brittle steel or ceramic)

Gradient often not in scale with elastomer material, e.g. $E_0 << E$.

Typical strains are just 0.1…0.5 %.
Hyperelastic material (3)

Elastomer material in comparison with metals and plastics:

- Energy-elasticity: Loading changes the distance of the atoms within the lattice of the metal and so increases the internal energy. When unloading it, this energy is immediately set free, the initial shape appears again.

- Entropy-elasticity: Within an elastomer, its macromolecules are balled if unloaded. During loading, a stretching and unballing appears. After unloading, more or less the unordered state appears again.

- Viscous behavior: every loading leads to an even small remaining deformation (creeping, relaxation).
Material laws for hyperelastic materials (1)

- The nominal or engineering strain is defined as the change in length divided by the original length:
  \[ \varepsilon = \frac{l_1 - l_0}{l_0} = \frac{\Delta l}{l_0} \]

- The stretch ratio \( \lambda \) now is another fundamental quantity to describe material deformation. It is defined as the current length divided by the original length:
  \[ \lambda = \frac{l_1}{l_0} = \frac{l_1 - l_0 + l_0}{l_0} = \varepsilon + 1 \]

- Analog to the three principal strains, we obtain from the principal axis transformation the three principal stretch ratios \( \lambda_1, \lambda_2, \lambda_3 \).

- The three stretch invariants (because independent from the used coordinate system) of the characteristic equation are analog:
  \[ I_1 = \lambda_1^2 + \lambda_2^2 + \lambda_3^2 \]
  \[ I_2 = \lambda_1^2 \lambda_2^2 + \lambda_2^2 \lambda_3^2 + \lambda_3^2 \lambda_1^2 \]
  \[ I_3 = \lambda_1^2 \lambda_2^2 \lambda_3^2 = 1 + \left( \frac{\Delta V}{V} \right)^2 = J^2 \]

  with \( J \): total volumetric ratio; if incompressible = 1
Material laws for hyperelastic materials (2)

- The description of the strain energy density $W$ is much more complex compared to linear elastic material, where the stress is just a linear function of strain.
- For hyperelastic material, the second Piola-Kirchhoff stress*) is defined from strain energy density function and Green-Lagrange strain (first derivative).
- In general, the strain energy density function in hyperelastic material is a function of the stretch invariants $W = f(I_1, I_2, I_3)$ or principal stretch ratios $W = f(\lambda_1, \lambda_2, \lambda_3)$, which is described in more detail on the next slides.

Constraints on the strain energy function $W$:

a) Zero strain = Zero energy: $W(0)=0$ (no energy is stored, if not loaded)
b) Zero strain = Zero stress: $W'(0)=0$ (unloaded condition)
c) Second derivative must be positive: $W''(\varepsilon) = \sigma'(\varepsilon) > 0$ for all $\varepsilon$ (stress always increases if strain increases, otherwise instability!)

*) Whereas the 1st Piola-Kirchhoff stress relates forces in the current configuration to areas in the reference configuration, the 2nd Piola-Kirchhoff stress tensor relates forces in the reference configuration to areas in the reference configuration.
Material laws for hyperelastic materials (3)

- Because of the material incompressibility, the deviatoric (subscript d or with ‘bar’) and volumetric (subscript V) terms of the strain energy function are split. As a result, the volumetric term is a function of the volume ratio \( J \) only (remember \( J^2 = I_3 \)):

\[
W = W_d \left( \bar{I}_1, \bar{I}_2 \right) + W_v (J) \quad \text{or} \quad W = W_d \left( \bar{\lambda}_1, \bar{\lambda}_2, \bar{\lambda}_3 \right) + W_v (J)
\]

- So, \( W_d \) is the strain energy necessary to change the shape, \( W_v \) the strain energy to change the volume.

- For typical hyperelastic material models, often phenomenological models are used, where the strain energy function has the form:

\[
W = \sum_{i+j=1}^{N} C_{ij} \left( I_1 - 3 \right)^i \left( I_2 - 3 \right)^j + \sum_{k=1}^{N} \frac{1}{D_k} (J - 1)^{2k}
\]

The \( C_{ij} \) and \( D_k \) are material constants which have to be determined by tests.

- This means, the strain energy function is a polynomial function. Depending on its order, no (=single curvature), one or more inflection points in the stress-strain curve may appear. For the higher order functions, enough test data has to be supplied!
Material laws for hyperelastic materials (4)

- As mentioned, typical hyperelastic material models have the form:
  
  \[ W = \sum_{i+j=1}^{N} C_{ij} (I_1 - 3)^i (I_2 - 3)^j + \sum_{k=1}^{N} \frac{1}{D_k} (J - 1)^{2k} \]

- Mechanica now supports five hyperelastic material laws of such a type:
  
  - Neo-Hookean (is the most simple approach):
    \[ w = c_{10} (I_1 - 3) + \frac{1}{D_1} (J - 1)^2 \]
  - Mooney-Rivlin:
    \[ w = c_{10} (I_1 - 3) + c_{01} (I_2 - 3) + \frac{1}{D_1} (J - 1)^2 \]
  - Polynomial form of order 2:
    \[ w = c_{10} (I_1 - 3) + c_{01} (I_1 - 3) + \frac{1}{D_1} (J - 1)^2 + c_{20} (I_1 - 3)^2 + c_{11} (I_1 - 3) (I_2 - 3) \]
  - Reduced Polynomial form of order 2:
    \[ w = c_{10} (I_1 - 3) + c_{20} (I_1 - 3)^2 + \frac{1}{D_1} (J - 1)^2 + \frac{1}{D_2} (J - 1)^4 \]
  - Yeoh (proposed not to use the second invariant term \( I_2 \), since it is more difficult to measure and provides less accurate fit for limited test data):
    \[ w = c_{10} (I_1 - 3) + c_{20} (I_1 - 3)^2 + c_{30} (I_1 - 3)^3 + \frac{1}{D_1} (J - 1)^2 + \frac{1}{D_2} (J - 1)^4 + \frac{1}{D_3} (J - 1)^6 \]
Material laws for hyperelastic materials (5)

- The sixth material law supported by Mechanica, Arruda-Boyce, has a slightly different form (it is not a phenomenological, but a micromechanical model):

\[ W = \mu \left\{ \frac{1}{2}(\bar{I}_1 - 3) + \frac{1}{20}\lambda_m^{-2}(\bar{I}_1^2 - 9) + \frac{11}{1050}\lambda_m^{-4}(\bar{I}_1^3 - 27) + \frac{19}{7000}\lambda_m^{-6}(\bar{I}_1^4 - 81) + \frac{519}{673750}\lambda_m^{-8}(\bar{I}_1^5 - 243) \right\} + \frac{1}{D} \left( \frac{J_e^2 - 1}{2} - \ln J_e \right) \]

- Here, the material constants have physical meaning: \( \mu = G_0 \) as initial shear modulus, \( \lambda_m \) as the limiting network stretch and \( D = 2/K_0 \) as the incompressibility parameter. This model is based on statistical mechanics; the coefficients are predefined functions of the limiting network stretch \( \lambda_m \). This is the stretch in the stretch-strain curve at which stress starts to increase without limit. If \( \lambda_m \) becomes infinite, the Arruda-Boyce form becomes the Neo-Hookean form!

- Remark for all material models: \( J_e \) is just the elastic volume ratio given by

\[ J_e = \frac{J}{J_{th}} = \frac{J}{(1 + \varepsilon_{th})^3} \]

with \( J = \) the total volumetric ratio, \( J_{th} = \) thermal volume ratio

Material laws for hyperelastic materials (6)

Some general remarks:

- The initial shear and initial bulk modulus, $G_0 = E_0/(2(1+\nu))$ and $K_0 = E_0/(3(1-2\nu))$, can be described with help of the material constants, for example in the material models of Neo-Hookean and Yeoh:
  
  $G_0 = 2C_{10}$
  
  $K_0 = \frac{2}{D_1}$

- For Mooney-Rivlin for example, the initial shear modulus becomes:
  
  $G_0 = 2(C_{10} + C_{01})$

- Which is the equivalent Poisson ratio used?

  The Poisson ratio used in the analysis can be determined from the used values for the initial shear and initial bulk modulus by the equation

  $$\nu = \frac{3K_0 - 2G_0}{6K_0 + 2G_0}$$

  For example, if $K_0/G_0 = 1000$, $\nu \approx 0.4995$
About selecting the material model and performing tests (1)

What is the “right” model to describe my material?

- If the strain is below approx. 5-10 %, for many applications the simple Hooke’s law is accurate enough to describe hyperelastic materials, so the time-consuming nonlinear analysis can be replaced by a very quick linear one.

- If the strain becomes bigger, but no or not a lot of test data is available, it is a good idea to start as a rough estimate with the most simple model, Neo Hookean:
  - In a first step, incompressibility can be assumed by setting $\nu=0.5$ or close to 0.5.

- In the literature, some (rough) empirical formulas can be found for the relation of the Shore-hardness $H$ and the shear modulus $G_0$ or initial E-modulus $E_0$; for example:
  - Battermann & Köhler: $G_0 = 0.086*1.045^H$
  - Rigbi ($H=$Shore A hardness): $H = 35.22735 + 18.75847 \ln(E_0)$

- Finally, the only two necessary material constants $C_{10}=G_0/2$ and $D_1=2/K_0$ can be simply obtained from the initial shear and initial bulk modulus, $G_0 = E_0/(2(1+\nu))$ and $K_0 = E_0/(3(1-2\nu))$ or $K_0=2G_0/(1+\nu)/(3(1-2\nu))$, like shown on the previous slide. If $\nu=0.5$, then we of course have $K_0=\infty$ and so $D_1=0$.

- If more test data is available, it is possible to let Mechanica select the best suitable material model. However, be very careful when the analysis is done for strains bigger than the maximum strain measured in the test! The higher-order material models in this case do not necessarily provide a higher accuracy!
About selecting the material model and performing tests (2)

How do I have to derive the right characteristic curve from the test?

- First of all, it is important that the strain rate applied in the test should be as close as possible to the strain rate applied in the later application, so that an accurate analysis can be performed for exactly this state!

- Elastomer material typically shows a hysteresis and an effect called “stress softening” *): After some cycles, the stress related to a certain strain decreases. This effect is not taken into account by the models previously described, so you have to perform the following treatment like shown in this sketch:

  - Select that cycle from the test data set for which you want to analyze your model:
    - loading or unloading
    - initial or n\textsuperscript{th} cycle
  - Subtract offset strain and stress
  - Perform curve fitting

*) For a possible model describing this effect, look for example in:
H.J. Qi, M.C. Boyce (Department of Mechanical Engineering, Massachusetts Institute of Technology, Cambridge, MA 02139, USA): Constitutive model for stretch-induced softening of the stress–stretch behavior of elastomeric materials; Journal of the Mechanics and Physics of Solids, Received 1 December 2003; accepted 14 April 2004
Implementation of hyperelastic material laws in Mechanica (1)

General remarks:

- Mechanica uses p-order finite element implementation to analyze hyperelastic materials. One of the advantages is that no special procedure is needed when the Poisson's ratio approaches to 0.5

- Literature sources for high-order finite element method can be found in the book “Finite Element Analysis” by Barna Szabo and Ivo Babuska. Specifically, on page 188, “In p-extension, the rate of converge (energy norm) is not affected by Poisson's ratio.” And on page 209, “Hence, locking does not occur. The elements can deform while preserving constant volume.”

Handling of nearly incompressible material in Mechanica:

If the value specified for \( D_1 = 2/K_0 \) is less than \( 1/500G_0 \), Mechanica uses this value as limit for \( D_1 \). So, we obtain for the maximum possible Poisson’s ratio used to “approximate ideal incompressibility”:

\[
K_0 = \frac{2}{D_1} = 1000G_0 \Rightarrow \frac{E_0}{3(1−2\nu)} = \frac{1000E_0}{2(1+\nu)} \Rightarrow \nu = \frac{1499}{3001} \approx 0.4995
\]

Remark: In linear elastic material analysis, the max. possible Poisson ratio Mechanica supports is 0.4999
Implementation of hyperelastic material laws in Mechanica (2)

Supported model/element types for hyperelastic material analysis:

- Large displacement analysis (LDA) is required for hyperelastic material analysis. All model/element types that support LDA also support hyperelastic material:
  - 3D volumes
  - 2D plane stress
  - 2D plane strain
  - 2D axial symmetry (will be supported in Wildfire 6)
  - Actually no support of beams and shells

- LDA: The forces and moments are equated iteratively at the deformed structure, as opposed to to SDA (small displacement analysis). Hence, an iterative procedure must be used to solve the nonlinear matrix equation for static analysis $K(u,f)u=f$

- Mechanica uses a modified Newton-Raphson procedure for this. To increase speed, BFGS (Broyden–Fletcher–Goldfarb–Shanno method) is used so that the stiffness matrix does not have to be computed and decomposed as often. A line search technique is used to control step size (reference: Bather, Klaus-Jürgen, Finite Element Procedures in Engineering Analysis, Prentice-Hall 1982)
Implementation of hyperelastic material laws in Mechanica (3)

Achieving convergence of the nonlinear matrix equation $K(u,f)\cdot u=f$ using Newton-Raphson technique:

- Before convergence we can calculate the residual error corresponding to the latest solution of the displacement vector $u$: $r=f-Ku$. Here, the residual vector $r$, has the dimensions of force (this force must be zero for system convergence). The Newton-Raphson solution then solves for $Kd\mathbf{u}=r$ to determine the change in $\mathbf{u}$ in the next iteration.

- The residual norm is the dot product $r \cdot d\mathbf{u}$. It can be thought of physically as a residual energy, which should be zero when we’re converged. We normalize the residual norm with the dot product of the total displacement and the total force vector, so the residual norm is: $(r \cdot d\mathbf{u})/(\mathbf{u} \cdot f)$.

- This residual norm must be smaller than the default value of 1.0E-14 to achieve convergence for the "Residual Norm Tolerance" in Mechanica (see .pas-file)

Defining hyperelastic material parameters in Mechanica (1)

The user has the following three options to define hyperelasticity:

- Select one of the 6 implemented material models and enter the necessary material constants manually
- Enter test data. Mechanica uses a Least Square Fitting algorithm (minimizing the normalized stress errors) to calculate the constants from the input test data for each material model. Then select the material model manually
- Let Mechanica automatically choose the material model with the best fit in the test domain based on the Root Mean Stress error

Check of the different material models from test data input:

- Mechanica performs a check on the stability of the material for six different forms of loading for $0.1 \leq \lambda \leq 10.0$ in intervals of $\Delta \lambda = 0.01$. The forms of loading are:
  - Uniaxial tension and compression
  - Equibiaxial tension and compression
  - Planar tension and compression
- For each loading type and $\lambda$, the tangential stiffness $D$ must be $> 0$
Defining hyperelastic material parameters in Mechanica (2)

Treating material model instability

- If an instability is found, Mechanica marks the model in the test data form with an exclamation mark and will not select it automatically (even though it may have a very small RMS error!)

- If the user overrides this by manually selecting the instable material model, Mechanica issues a warning message with the values of $\varepsilon$ for which instability is observed ($-0.9 \leq \varepsilon \leq 9.0$)

- The model may be used just up to these limits, otherwise the analysis will fail!

The following four types of tests are supported:

- Uniaxial: Uniaxial tension
- Biaxial: Equibiaxial tension
- Planar: A certain plain strain condition (as described later)
- Volumetric: Hydrostatic pressure

These stress/strain & stretch states are depicted on the next slide, respectively.
Defining hyperelastic material parameters in Mechanica (3)

Idealized stress/strain states of the four hyperelastic material tests supported in Mechanica:

- **Uniaxial:**
  \[ \lambda_1 = \frac{1}{\lambda_2^2} = \frac{1}{\lambda_3^2} \]

- **Planar:**
  \[ \lambda_1 = \frac{1}{\lambda_2}, \lambda_3 = 1 \]
  \( (\sigma_3 \text{ is positive because of lateral strain suppression in 3-direction, but not applied as external force like } \sigma_1) \)

- **Equibaxial:**
  \[ \lambda_1 = \lambda_2 = \frac{1}{\sqrt{\lambda_3}} \]

- **Volumetric:**
  \[ \lambda_1 = \lambda_2 = \lambda_3 = J^{1/3} \leq 1 \]
  \( (\sigma_3 \text{ (hydrostatic pressure } p=\sigma_1=\sigma_2 =\sigma_3) \)

**Remark:**
Stresses (red) or strains (blue) where no arrow is shown are Zero!
Defining hyperelastic material parameters in Mechanica (4)

Obtaining the material constants $C_{ij}$ and $D_k$ from the test data input:

- From the uniaxial, equal biaxial, and planar tests, only the $C_{ij}$ are determined. The material is assumed to be incompressible, if no additional volumetric test data is given! In this case, the $D_k$'s are shown as 0 (meaning incompressible). Remember, the engine assumes a *nearly* incompressible material then and uses $D_1 = 1/(500 \, G_0)$ during the analysis like previously described (means $\nu=0.4995$)

- From the volumetric test, only the $D_k$'s are being estimated; in this case of course, incompressibility is not assumed. The $C_{ij}$ cannot be calculated from this test because hydrostatic pressure just creates a volume change and no shape change! That’s why a volumetric test alone is not sufficient to characterize hyperelastic material

- If more than one test is entered, then the data from all of the tests are considered when determining the material properties. No one test counts more than any of the others; all tests are considered equally

Important Remark:
In general, engineering (nominal) values have to be entered for stress and strain into the test data forms!
Test set-ups and specimen shapes of the supported material tests (1)

Uniaxial:

- This is the classical uniaxial tension rod mounted into a tensile testing machine

Note: The strain must of course be measured in the thinner area of the test rod, for example by optical scanning (video extensometry); the thicker parts of the tension rod which are clamped must not be taken into account!

Example taken from paper CMMT(MN)054:

“Test Methods for Determining Hyperelastic Properties of Flexible Adhesives”

Bruce Duncan (1999)
Centre for Materials Measurement and Technology
National Physical Laboratory
Queens Road, Teddington, Middlesex, TW11 0LW
Telephone: 020 8977 3222 (switchboard)
Direct Line: 020 8943 6795
Facsimile: 020 8943 6046
E-mail: bruce.duncan@npl.co.uk
Test set-ups and specimen shapes of the supported material tests (2)

Biaxial:

- This is a disk under equibiaxial tension. The specimen mounted into a “scissor” fixture for an uniaxial testing machine and the stress state may look as follows:

For this specimen type, failure will occur in the edges where the load is introduced.

Reference: CMMT(MN)054
Test set-ups and specimen shapes of the supported material tests (3)

Biaxial (cont’d):

- Another test setup and specimen for equibiaxial tension may look like this:

Reference: www.axelproducts.com
Test set-ups and specimen shapes of the supported material tests (4)

Planar:

- A thin sheet of hyperelastic material is clamped, so that lateral strains are prohibited here, and pulled!

Acc. to ref. CMMT(MN)054, the planar test results shall be relatively insensitive to the grip separation “d”, but this should be treated with care for larger strains. See the planar test example in part 2 of this presentation!
Test set-ups and specimen shapes of the supported material tests (5)

Volumetric

- A volumetric test setup like this compresses a cylindrical elastomer specimen constrained in a stiff fixture.

- The actual displacement during compression is very small and great care must be taken to measure only the specimen compliance and not the stiffness of the instrument itself.

- The initial slope of the resulting stress-strain function is the bulk modulus. This value is typically 2-3 orders of magnitude greater than the shear modulus for dense elastomers.

Reference: www.axelproducts.com
The uniaxial compression test (1)

Simple compression:

- Biggest problem of this test is that lateral strains are disturbed by friction effects.
- From the analysis results shown below, one can conclude that even very small levels of friction significantly affect the measured stiffness. Furthermore, this effect is apparent at both low and high strains. This is particularly troubling because friction values for elastomers are typically a function of normal force and are not well characterized.
- As such, the experimental compression data cannot be corrected with a significant degree of certainty.
- Unfortunately, both tension and compression information is valuable to obtain because unlike some metal material models, elastomers behave very differently in compression than in tension!

Reference: www.axelproducts.com
The uniaxial compression test (2)

According to Axel Testing Services, the equibiaxial extension experiment also provides compression information:

- As an elastomer is radially strained in all directions in a single plane, the free surfaces come together.

- For incompressible materials, the state of strain in the material is the same as that in simple compression (if free from friction!). The measured experimental parameters are radial strain and stress.

- These biaxial strains and biaxial stresses can be converted directly to compression strains and compression stresses as follows:

\[
\sigma_c = \sigma_b \left(1 + \varepsilon_b\right)^3
\]

\[
\varepsilon_c = \frac{1}{(\varepsilon_b + 1)^2} - 1
\]

- \(\sigma_c\): nominal compression stress
- \(\sigma_b\): nominal biaxial extension stress
- \(\varepsilon_c\): nominal compression strain
- \(\varepsilon_b\): nominal biaxial extension strain

- It typically isn’t necessary to do this conversion because most curve fitters accept equibiaxial extension data directly (like Mechanica).
Stress and strain definitions in the Mechanica LDA analysis (1)

Material test definitions:
- Remember: In Mechanica, engineering (nominal) values have to be entered for stress and strain into the hyperelastic test data forms!

Analysis results:
- Mechanica reports true stresses in the LDA analysis (in SDA of course, nominal values are output: There is no significant difference between true and nominal stress!)
- Unlike for the test definitions, Mechanica does not use engineering strains in LDA. Mechanica reports the so called “Eulerian” or “Almansi” strain, which becomes surprisingly small for large nominal tension strains and very big for large negative strains (theoretic maximum for infinite nominal tension strain is just 0.5!)
- The reason is that this strain is defined with respect to the current configuration (stretched length $l_1$) of the body – not the initial length $l_0$!
- For further explanation, the next slide shows the equations for the different strain definitions
Stress and strain definitions in the Mechanica LDA analysis (2)

Strain definitions:


- The various strain measures have the following values for a tensile rod ($l_0$ is the initial length, $l_1$ is the current length):
  - Infinitesimal, “engineering” or Cauchy strain: (for small displacement problems only)
    \[ \varepsilon = \frac{l_1 - l_0}{l_0} = \frac{\Delta l}{l_0} \]
  - Logarithmic (“natural”, “true”, “Hencky”) strain: (obtained by integrating the incremental strain)
    \[ \partial \varepsilon_L = \frac{\partial l}{l_1} \Rightarrow \int \partial \varepsilon_L = \int_{l_0}^{l_1} \frac{\partial l}{l_1} \Rightarrow \varepsilon_L = \ln \left( \frac{l_1}{l_0} \right) = \ln \lambda \]
    \[ \Leftrightarrow \varepsilon_L = \ln(1 + \varepsilon) = \varepsilon - \frac{\varepsilon^2}{2} + \frac{\varepsilon^3}{3} - \frac{\varepsilon^4}{4} + \ldots \]
  - Green-Lagrange Strain: (defined with respect to the initial configuration)
    \[ \varepsilon_G = \frac{1}{2} \frac{l_1^2 - l_0^2}{l_0^2} \]
  - Eulerian (Almansi) Strain: (defined with respect to the deformed configuration)
    \[ \varepsilon_E = \frac{1}{2} \frac{l_1^2 - l_0^2}{l_1^2} \]


- The various strain measures have the following values for a tensile rod ($l_0$ is the initial length, $l_1$ is the current length):
  - Infinitesimal, “engineering” or Cauchy strain: (for small displacement problems only)
    \[ \varepsilon = \frac{l_1 - l_0}{l_0} = \frac{\Delta l}{l_0} \]
  - Logarithmic (“natural”, “true”, “Hencky”) strain: (obtained by integrating the incremental strain)
    \[ \partial \varepsilon_L = \frac{\partial l}{l_1} \Rightarrow \int \partial \varepsilon_L = \int_{l_0}^{l_1} \frac{\partial l}{l_1} \Rightarrow \varepsilon_L = \ln \left( \frac{l_1}{l_0} \right) = \ln \lambda \]
    \[ \Leftrightarrow \varepsilon_L = \ln(1 + \varepsilon) = \varepsilon - \frac{\varepsilon^2}{2} + \frac{\varepsilon^3}{3} - \frac{\varepsilon^4}{4} + \ldots \]
  - Green-Lagrange Strain: (defined with respect to the initial configuration)
    \[ \varepsilon_G = \frac{1}{2} \frac{l_1^2 - l_0^2}{l_0^2} \]
  - Eulerian (Almansi) Strain: (defined with respect to the deformed configuration)
    \[ \varepsilon_E = \frac{1}{2} \frac{l_1^2 - l_0^2}{l_1^2} \]
Stress and strain definitions in the Mechanica LDA analysis (3)

Graphical representation of the different strains:

\[ \varepsilon_G = \frac{1}{2} \frac{l_1^2 - l_0^2}{l_0^2} \]

\[ \varepsilon = \frac{l_1 - l_0}{l_0} = \lambda - 1 \]

\[ \varepsilon_L = \ln \left( \frac{l_1}{l_0} \right) = \ln \lambda \]

\[ \varepsilon_E = \frac{1}{2} \frac{l_1^2 - l_0^2}{l_1^2} \]

Reported by Mechanica in LDA! stretch \( \lambda = \frac{l_1}{l_0} \)
Part 2

Application examples
A test specimen subjected to uniaxial load (1)

Test specimen and test data:

- Provided uniaxial test data (engineering values) of an example elastomer:

<table>
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<th>$\varepsilon$ [%]</th>
<th>$\sigma_{\text{tech}}$ [MPa]</th>
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<td>5.922</td>
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<tr>
<td>6.449</td>
<td>9.046</td>
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</table>

Test specimen acc. to DIN 53504-S2

Loaded cross section is 8 mm²

Measurement length 20 mm

Clamped length 50 mm

R12.5

Schnitt 0
A test specimen subjected to uniaxial load (2)

For the Arruda-Boyce model, the Least Square Fitting algorithm failed, so it cannot be used (and is not displayed). The exclamation mark means, that for a certain strain range, the model is unstable (Zero tangent stiffness).

Mechanica automatically selects “Polynomial Order 2” – model as best fit to test data. The “D”-values are shown as zero (≈incompressible), since no volumetric test has been specified. However, internally Mechanica uses \( D_1 = D_2 = 1/(500G) \), which corresponds to a Poisson ratio of 0.4995.

If this box is unchecked, you can manually enter the coefficients for the material law and compare them with the test points in the graph.

- \( G_0 = 2(C_{10}+C_{01}) = 3.11234 \text{ MPa} \)
- \( K_0 = 2/D_1 = 3112.34 \text{ MPa (v=0.4995)} \)
- \( E_0 = 2G_0(1+v) = 9.3339077 \text{ MPa} \)
A test specimen subjected to uniaxial load (3)

Model Set-Up

- There are different ways to set up the FEM-model of the specimen:
  - The most “realistic” one is to use the complete specimen geometry from Pro/E, prepare and mesh it, define the material and run the analysis. This looks of course nicest
  - However, to save time we will just run an eighth part of the measurement-zone of the model with symmetry constraints, containing just three bricks (you could also use 2D plane stress of course!). This analysis will run very fast!
- We create some measures to determine the engineering values for stress and strain, since the Mechanica engine reports only true stress and Almansi strain:
  - A measure for nominal strain, using the following formula that derives the engineering strain $\varepsilon$ from the Eulerian (Almansi) strain $\varepsilon_E$ output by Mechanica:
    \[
    \varepsilon_E = \frac{1}{2} \frac{l_i^2 - l_0^2}{l_i^2} = \frac{1}{2} \left( 1 - \frac{1}{\varepsilon + 1} \right)^2 \Rightarrow \varepsilon = \sqrt{\frac{1}{1 - 2\varepsilon_E} - 1}
    \]
  - As cross-check, a measure for nominal strain, derived from the specimen length change divided by the initial length (with help of a Mechanica computed measure)
  - A computed measure for nominal stress, using the constraint reaction force divided by the initial cross section
A test specimen subjected to uniaxial load (4)

FEM model

- Full model:

- Eighth part of the measurement-zone of the model with symmetry constraints:
A test specimen subjected to uniaxial load (5)

Results (displacement magnitude, deformed shape in scale)

- Full model:

- Eighth part of the measurement-zone of the model with symmetry constraints:

- Eighth part just with linear behavior (small strain solution and properties):
A test specimen subjected to uniaxial load (6)

Stress-strain-curves of the example elastomer

Here, a good match between the linear and hyperelastic material model is just prevailing below strains of approx. 5%!
A test specimen subjected to uniaxial load (7)

Engineering stress and strain versus stretch

Engineering strain (FEM LDA)
- Almansi strain (FEM LDA)
- Engineering strain (FEM linear)

Almansi strain goes for theoretical limit of 0.5 (0.491 for max. load applied)
A volumetric compression test (1)

Test specimen and test data:
- Like shown in the section about material tests, in a volumetric test a cylindrical specimen is uniaxial compressed while it is constrained in a very stiff fixture.
- From Hooke’s law, we have with $\sigma_1 = \sigma_{ax} = F/A$, $\sigma_2 = \sigma_3 = \sigma_q$ and $\varepsilon_2 = \varepsilon_3 = 0$:
  
  $\varepsilon_1 = \frac{1}{E} \left( \sigma_1 - \nu (\sigma_2 + \sigma_3) \right) = \frac{1}{E} \left( \frac{F}{A} - 2\nu \sigma_q \right) = \varepsilon_{ax}$

  $\varepsilon_2 = \frac{1}{E} \left( \sigma_2 - \nu (\sigma_1 + \sigma_3) \right) = \frac{1}{E} \left( \sigma_q - \nu \left( \frac{F}{A} + \sigma_q \right) \right) = \varepsilon_q = 0$

  $\varepsilon_3 = \frac{1}{E} \left( \sigma_3 - \nu (\sigma_1 + \sigma_2) \right) = \frac{1}{E} \left( \sigma_q - \nu \left( \frac{F}{A} + \sigma_q \right) \right) = \varepsilon_q = 0$

  So, we have two different equations to solve for the two unknown quantities $\sigma_q$ and $\varepsilon_{ax}$. We obtain:

  $\sigma_q = \frac{F}{A} \frac{\nu}{1-\nu}$

  $\varepsilon_{ax} = \frac{1}{E} \frac{F}{A} \left( 1 - 2 \frac{\nu^2}{1-\nu} \right)$
A volumetric compression test (2)

Analytical calculation of volumetric test specimen behavior:

- Our cylindrical specimen has a diameter of 5 mm and a length of 20 mm. We apply a force of \( F = 100 \) N. We use the same elastomer material like in the uniaxial test before. We obtain with \( \nu = 0.4995 \) and \( E = 9.3339077 \) MPa:

  \[
  A = \frac{\pi d^2}{4} = 19.63495 \text{ mm}^2
  \]

  \[
  \sigma_{\text{ax}} = \frac{F}{A} = 5.092958 \text{ MPa}
  \]

  \[
  \sigma_{\text{q}} = \frac{F}{A} \frac{\nu}{1-\nu} = 5.082782 \text{ MPa}
  \]

  \[
  \varepsilon_{\text{ax}} = \frac{1}{E} \frac{F}{A} \left( 1 - 2 \frac{\nu^2}{1-\nu} \right) = 0.00163474
  \]

  \[
  \Delta l = \varepsilon_{\text{ax}} l = 0.0326948 \text{ mm} \approx 33 \mu m
  \]

- With \( \nu = 0.4999 \), the maximum value supported in Mechanica for linear materials, we obtain:

  \[
  \sigma_{\text{q}} = 5.090921 \text{ MPa}, \quad \varepsilon_{\text{ax}} = 0.000327297 \text{ and } \Delta l = 6.546 \mu m
  \]
A volumetric compression test (3)

Comparison of hyperelastic material with steel:

- If we compress a steel cylinder with the same dimensions, we obtain with $E=210000$ MPa for the unconstrained condition ($\sigma_q=0$):
  
  $K = \frac{EA}{l}; \quad F = K\Delta l$

  $\Delta l = \frac{Fl}{EA} = \frac{4Fl}{E\pi l^2} = 0.485 \mu m$

- With $\nu=0.3$ we obtain (constrained condition $\varepsilon_q=0$):
  
  $\sigma_q=2.183$ MPa, $\varepsilon_{ax}=1.8016 \times 10^{-5}$ and $\Delta l=0.36 \mu m$

Because of the compressibility of steel, there is not a big difference to the unconstrained condition (factor $\approx 1.35$)!

- The elastomer cylinder, assuming $\nu=0.4999$, deforms just $18$ times more ($\Delta l=6.546 \mu m$) than the steel cylinder, even though its E-modulus is approx. $22500$ times lower!

→ Elastomer material can behave surprisingly stiff under certain conditions! Take this into account when designing for example “soft” rubber layers to homogenize bearing stress!
A volumetric compression test (4)

Model Set-Up:

- Since we have for this condition just very small deformations, we could run the compression test analysis with the linear theory as 2D axial symmetric model (in WF6, 2D axial symmetric models will also support LDA and hyperelasticity).

- Mechanica in this case automatically selects the initial values $G_0$ and $K_0$ from the example elastomer test data input (equivalent to $E_0=9.3339$ and $\nu=0.4995$).

- However, we will run the compression test as 2D plane strain model. This is possible, since for this loading condition the axial displacement for a given axial stress is not a function of the specimen cross section, but just of its length.

- In 2D plane strain, we can run the model linearized and with LDA including hyperelasticity, to check the difference!
A volumetric compression test (5)

Model Set-Up:

- For this simple stress state without any gradients, a very coarse mesh is sufficient to obtain accurate results:

- Also in this model, we define a computed measure for the engineering strain to compare it with the Almansi strain reported in LDA. The difference should be negligible here!
A volumetric compression test (6)

Analysis results:

- As expected, the difference between analytical results, SDA with linear material data and LDA with hyperelastic material is very small for volumetric compression:

  **Linear analysis results (Multi Pass):**
  - max_stress_xx: \(-5.082782e+00\) 0.0%
  - max_stress_yy: \(-5.092958e+00\) 0.0%
  - max_stress_zz: \(-5.082782e+00\) 0.0%
  - strain_energy: \(2.081421e-01\) 0.0%
  - displacement_Y: \(-3.269489e-02\) 0.0%
  - epsilon_ax_eng: \(-1.634744e-03\) 0.0%

  **Linear analysis results (analytical solution):**
  - \(\sigma_{xx} = \frac{F}{A} = 5.092958\) MPa
  - \(\sigma_y = \frac{F}{A} \frac{\nu}{1-\nu} = 5.082782\) MPa
  - \(\varepsilon_{ax} = \frac{1}{E} \left(1-2\nu\right) \left(1-\frac{\nu^2}{1-\nu}\right) = 0.00163474\)
  - \(\Delta l = \varepsilon_{ax} l = 0.0326948\) mm

  **LDA results with hyperelasticity (Single Pass):**
  - max_stress_xx: \(-5.082742e+00\)
  - max_stress_yy: \(-5.092958e+00\)
  - max_stress_zz: \(-5.082742e+00\)
  - strain_energy: \(2.077312e-01\)
  - Almansi_strain: \(-1.638205e-03\)
  - displacement_Y: \(-3.269489e-02\)
  - epsilon_ax_eng: \(-1.634191e-03\)

No significant difference!
A planar test (1)

Specimen Geometry:

- We will use a thin sheet of the same example elastomer: Thickness 2 mm, clamped length 100 mm, grip separation d=30 mm

lateral strain $\varepsilon_3=0$ because of clamping!
A planar test (2)

Test specimen idealization:

- The most simple way to analyze this specimen is a 2D plane stress model (even though it contains a planar strain state – theoretically at least in the measurement zone in the center)
- Quarter symmetry can be used for this plane stress idealization (shown in red)
- The central, vertical cross section of the specimen (shown in blue) could also be idealized as plane strain condition (just if the strain state there is sufficiently planar, what we will subsequently examine!)
- However, this idealization does not allow to check for aberrations from the ideal plane strain state at the outer borders of the specimen, which will influence the necessary tension force and so the engineering stress
- Running this model with 3D solids does not give any advantage!
A planar test (3)

Model Setup:

- 2D plane stress, quarter symmetry
- Several measures have been created to track the nonlinear behavior especially in the center of the measurement zone

An enforced displacement of 30 mm, applied in increments, creates a maximum nominal strain of 200 % \( (\lambda_y = 3) \)

\[
\varepsilon_E = \frac{1}{2} \left( \frac{l_1^2 - l_0^2}{l_1^2} \right) = \frac{1}{2} \left( 1 - \left( \frac{1}{\varepsilon + 1} \right)^2 \right) \iff \varepsilon = \sqrt{\frac{1}{1 - 2\varepsilon_E}} - 1
\]
A planar test (4)

Displacement results (in scale):

Enforced displacement: 3 mm

Enforced displacement: 6 mm

Enforced displacement: 18 mm

Enforced displacement: 24 mm

Enforced displacement: 12 mm

Enforced displacement: 30 mm

Extreme lateral contraction!
A planar test (5)

True stress and Almansi strain results for 30 mm displacement ($\lambda_y=3$):

- **x-strain along curve 1** (this strain should be Zero for a plane strain condition!)
- **z-strain along curve 1** (out-of-plane-strain)
- **y-strain along curve 1**
- **y-strain along curve 2** (becomes smaller when reaching the clamping zone!)

Geometric center of the specimen
A planar test (6)

Strain results versus stretch in the geometric center of the specimen:

The graph shows engineering and Almansi strains for comparison.

These strains are the y-values for the geometric center of the specimen (in the measurement zone, see test setup description).

Here, the “nominal” engineering strain in y is the “averaged” strain from the total specimen elongation under y-load.
A planar test (7)

Engineering strain results vs. stretch in the specimen geometric center:

These graphs show engineering strains only, because this is more familiar for most users!

We have an approximate plane strain condition just under very small loads! (green line = x-strain should be Zero for all stretch values to have a true plane strain condition!)

For higher loadings, we obtain more and more a uniaxial stress & triaxial strain state in the measurement zone!
A planar test (8)

Conclusions:

- In nonlinear analysis of hyperelastic material, the stress and strain state quality (type) may vary significantly during the analysis, not only the quantity like in linear analysis with metals.

- In the example shown, at the beginning we have a plane strain and plane stress condition (plane conditions in different planes, respectively). When the stretch is increasing, we obtain in the center of the specimen, where we measure the strains, more and more a uniaxial stress and a triaxial strain state, which of course is something that we don’t want to have!

- Hence, great care must be taken when defining multiaxial test geometry and load levels!

- For a precise test evaluation, also FEM analyses are recommended to understand the specimen behavior. Mechanica can help you a lot here!
Influence of the Material Law (1)

Motivation:

- Until now, all example analyses were based on the same simple uniaxial tension test and the “polynomial order 2” hyperelastic material model
- For comparison, we will now also examine the influence of
  - the material law used: We will run some example analyses not only with the proposed (“automatic - best fit”) model, but also with that one on the list with the second smallest RMS-error;
  - the stress and strain state (uniaxial or planar test)

Goal:

- Obtain a “feeling” for how sensitive our analyses and predictions are against such changes
- Learn what to do to minimize errors
Influence of the Material Law (2)

Material law influence

- We select the best and second best material law for the test data fit (see RMS error) and re-run both uniaxial and planar test, respectively

Polynomial order 2 model delivers lowest RMS error and is therefore automatically proposed by Mechanica

The Yeoh model has to be selected manually
Influence of the Material Law (3)

Uniaxial Test Case

We run the uniaxial model again for the tension case. The differences are very low in this case, like expected!
**Influence of the Material Law (4)**

**Planar Test Case**

- We run the planar model again. The differences are unexpectedly very big! (see some extreme cases in red)

<table>
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<tr>
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<th>Measure results of Yeoh material model</th>
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</tr>
</tbody>
</table>
Influence of the Material Law (5)

Planar Test Case

The lateral contraction in the planar test is predicted completely different with the two different material models (results are in scale for a nominal stretch of $\lambda=3$)!
Influence of the Material Law (6)

Conclusions:

- The difference in the lateral contraction cannot be explained with a different bulk modulus (means different Poisson ratios): For both models, since no volumetric tests have been performed, so the $D_i$ are set to Zero (internally $D_0=1/500$ $G_0$, this means the same Poisson ratio of 0.4995 is used both models)

- The Yeoh model neglects the second stretch invariant in the strain energy density function, just the first one is used, which may explain the difference

$$I_1 = \lambda_1^2 + \lambda_2^2 + \lambda_3^2$$
$$I_2 = \lambda_1^2 \lambda_2^2 + \lambda_2^2 \lambda_3^2 + \lambda_1^2 \lambda_3^2 \neq 0$$

- With this example, it becomes clear that the test conditions (uniaxial, planar...) highly influence the usability of test data for the analysis of the real design: A simple tension test is often not enough for a good prediction of your real part behavior if this not loaded in simple tension as well!

- In general, do as many different tests as possible to characterize your material!

- Especially use test conditions as close as possible to the loading state of the part you want to design!
Thanks for your attention!
PTC Simulation Services Introduction

PTC Global Services provides services for our own simulation products:

- Pro/ENGINEER Mechanica as a FEA tool with p-method for structural mechanical, thermal and thermo-mechanical analysis
- Pro/ENGINEER MDX and MDO (Mechanism Design Extension and Mechanism Dynamics Option) for kinematic and dynamic multi-body simulations

The benefits are accomplished as following:

- Required calculations
- Development of the required analysis and optimization, working with the design team, directly on the working CAD data, including adoption of mechanical systems engineering tasks
- On-site simulation consulting → Software and calculation method knowledge transfer
- Simulation training and workshops from PTC University

The following slides show the newest examples of simulation project and education references. Numerous other references from other clients and to other simulation issues can be provided upon request.
PTC Global Services Examines the Dynamic Structural Behavior of an Opto-Mechanical Subsystem Prototype from Carl Zeiss Optronics with Pro/ENGINEER® Mechanica® Software

Carl Zeiss Optronics GmbH, a member of the Carl Zeiss Group located in Oberkochen, Germany, develops and produces high-precision and robust opto-electronic systems for observation and defense purposes. For such products exposed to intense loading, advanced system analysis with the Finite Element Method is an integral part of the product development.

**BUSINESS INITIATIVE**

- During the development of new observation systems, Zeiss Optronics performs analyses to study the behavior of the installed subsystems for assuring that the final product works accurately. For the subsystem shown to the right, PTC Global Services was charged with these examinations.

**SOLUTION**

- PTC built up the dynamic analysis model in Mechanica with the help of given Zeiss Pro/ENGINEER CAD model assembly data. Reasonable linearizations were developed for the guiding system and the preloaded drive mechanism. Their validity was strictly controlled during all subsequent analyses. Given sine sweep test data was compared with the dynamic frequency analysis results to assure an accurate mathematical model of the subsystem.

**RESULT**

- Good match of test and analysis result data helped to understand the dynamic characteristics of the subsystem.
- Points in the design leading to unwanted behavior could be identified and solutions were provided.

“The dynamic analysis study gave us a very good understanding of exactly what happens in our newly designed subassembly for moving lens elements. PTC’s responsible consultant for the project, Dr. Roland Jakel, also provided excellent ideas for helpful design modifications. With the obtained knowledge, we can now enhance the subsystem in a very early stage of the development, ensuring that it meets the requirements.”

Dr.-Ing. Thomas Meenken, Team Leader Simulation, Carl Zeiss Optronics GmbH
PTC University Further Educates Otto Bock HealthCare in Advanced Nonlinear Contact and Bolt Analysis with Pro/ENGINEER® Mechanica® Software

Otto Bock HealthCare is the leading supplier of innovative products for people with restricted mobility, and, as a recognized system provider of high-quality, technologically advanced products and services, it is also the global leader in orthopedic technology. The company was founded in Berlin in 1919, and is now led by Professor Hans Georg Näder, the third-generation managing shareholder. In addition to the core competency as the leading company in the Orthobionic® field, Bionicmobility® is an additional competency of Otto Bock. It combines mobility solutions such as high-quality lightweight and active wheelchairs, power wheelchairs, and products for pediatric rehabilitation and seating shell systems.

**BUSINESS INITIATIVE**
- The technologically advanced orthopedic products developed by Otto Bock require extensive Finite Element analyses to ensure proper function in service over their complete life span. For a more accurate solution of nonlinear problems like contact and fastener analyses, Otto Bock wanted to deepen the knowledge of their design engineers in this demanding topic.

**SOLUTION**
- PTC offered an on-site workshop for these nonlinear analysis themes with the opportunity for Otto Bock engineers to get their typical product analysis tasks exemplarily solved by the PTC course instructor.

**RESULT**
- The theoretic background knowledge provided with help of the Otto Bock product examples supports the engineers in applying the Mechanica FEM code correctly to their analysis tasks.

“"We listened well to the background information PTC provided in this course for frictionless and infinite friction contact theory in Mechanica as well as to the extensive explanations about behavior of fasteners. In addition, the example solutions provided help us a lot since we can apply all this directly to our new products under development.”

Ralf Allermann, Development / Design / Simulation, Otto Bock HealthCare GmbH
The Vaillant Group is an internationally operating heating, ventilation and air-conditioning technology concern based in Remscheid, Germany. As one of the world's market and technology leaders, the company develops and produces tailor-made products, systems and services for domestic comfort. The product portfolio ranges from efficient heating appliances based on customary fuels to system solutions for using regenerative energy sources. As Europe’s number one heating technology manufacturer, ‘thinking ahead’ is a culture which is embraced throughout their business. To ensure an excellent product quality and short development cycles, Vaillant uses modern CAE tools like CFD software or Mechanica as a Finite Element program.

**BUSINESS INITIATIVE**

- During the product lifecycle, design modifications are often adopted to decrease manufacturing costs while maintaining or increasing product quality. Also here, FEM is used to ensure the reliability of such changes. The Mechanica software knowledge of the Vaillant CAx application engineers was to be extended to advanced nonlinear simulation, so that very ambitious analysis tasks can be solved in-house without external support.

**SOLUTION**

- PTC offered an individual simulation workshop focusing on contact and bolt analysis theory. Original Vaillant product examples and CAD data sets were used for practical training examples.

**RESULT**

- Obtained knowledge in nonlinear contact and bolt FEM analysis
- Obtained FEM sample solutions for typical Vaillant analysis tasks, like for the screw fitting shown right, allowing own further studies

“Attending the advanced PTC training in Mechanica nonlinear contact and bolt analysis has enabled us to do these ambitious expert analyses in the future without external support. We value the excellent knowledge transfer and the sample solutions accurately provided on base of our own products and simulation tasks.”

Stefan Schweitzer-De Bortoli, CAx Application Engineer Simulation Tools, Vaillant GmbH

**Right image:** A state-of-the-art Vaillant heating system for domestic comfort

**Left images:** Screwed joint analysis of a pipe connection with non-regular geometry, performed with Mechanica in the customized workshop. A hexagonal spigot nut (1) connects the copper tube end (2) with the brass tube (3), a sealing (4) is used against leakage of the fluid. Such bolted connections cannot be analyzed analytically acc. to bolt analysis guidelines because of their geometry; therefore Mechanica allows an accurate FEM analysis.
Dictionary Technical English-German (1)

Most important terminology for German listeners:

- bearing stress – Auflagerspannung
- bulk modulus – Kompressionsmodul \( K = -\frac{\Delta \rho \cdot V}{\Delta V} = \frac{E}{3(1-2\nu)} \)
- coefficient of thermal expansion (CTE) – Wärmeausdehnungskoeffizient \( \alpha \)
- density – Dichte
- dot (scalar) product – Skalarprodukt
- hardness – Härte
- modulus of elasticity – Elastizitätsmodul \( E \)
- nominal (or engineering) strain – technische Dehnung \( \varepsilon = \frac{\Delta l}{l} \)
- nominal (or engineering) stress – technische Spannung \( \sigma = \frac{F}{A_0} \)
- poisson ratio – Querdehnzahl \( \nu \)
- principal axis transformation – Hauptachsentransformation
Dictionary English-German (2)

- shear modulus – Schubmodul \( G = \frac{E}{2(1+\nu)} \)
- strain – Dehnung \( \varepsilon \)
- strain energy density function – Dehnungsenergiedichte-Funktion \( W \)
  (volumenbezogen)
- stress softening – Entfestigung
- stretch – Streckung, Längung
- stretch invariants – Streckungsinvarianten \( I_1, I_2, I_3 \)
- stretch ratio – Streckungsverhältnis \( \lambda = \varepsilon + 1 \)
- tension strength – Zugfestigkeit
- volumetric ratio – relative Volumenänderung \( J = \Delta V / V \)
Informations about the Presenter

Roland Jakel

- Dipl.-Ing. for mechanical engineering (Technische Universität Clausthal)
- Ph.-D. in design and analysis of engineering ceramics (FEM-Analysis and subroutine programming with Marc/Mentat)
- 1996-2001 Employee at Dasa in Bremen (Daimler-Benz Aerospace, Product Division Space-Infrastructure, today EADS Astrium):
  - Structural simulation (FEM-Analysis with NASTRAN/PATRAN and Mechanica)
  - Project management for Ariane 5 Upper Stage „ESC-A“ Subsystems (Stage Damping System “SARO”, Inter Tank Structure)
- At the former DENC AG („Design ENgineering Consultants“) from 2001-2005 responsible for structural simulation services and education with the PTC simulation products (Mechanica, MDX, MDO, BMX)
- Since the DENC AG acquisition by PTC in 2005, Roland Jakel is responsible for the PTC simulation services within the Global Services Organization (GSO) for CER (Central Europe)